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A NEW FAMILY OF LIFE DISTRIBUTIONS

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#### SUMMARY

A new two parameter family of life length distributions is presented which is derived from a model for fatigue. This derivation follows from considerations of renewal theory for the number of cycles needed to force a fatigue crack extension to exceed a critical value. Some closure properties of this family are given and some comparisons made with other families such as the lognormal which have been previously used in fatigue studies.

## 1. INTRODUCTION

It is well known that for the amount of fatigue data which can usually be obtained almost any two dimensional parametric family of distributions can be made to fit reasonably well. In fact, in the region of central tendency the lognormal, the Weibull, the Gamma, etc., can all be fitted by parametric estimation and because of the relatively small sample sizes hardly any can be rejected by, say a Chi-square Goodness of Fit test. However, when it becomes a question of predicting the "safe life" say the one thousandth percentile, there is a wide discrepancy between these models.

For this reason a family of distributions which is obtained from considerations of the basic characteristics of the fatigue process should be more persuasive in its implications than any ad hoc family chosen for extraneous reasons. In this paper we derive, using some elementary renewal theory, a two parameter family of nonnegative random variables as an idealization of the number of cycles necessary to force a fatigue crack to grow to a critical value. We then examine some of its relevant properties.

## 2. A MODEL FOR THE DISTRIBUTION OF LIFE

We propose a basic framework and notation which is similar to that used previously in [3]. We consider only standardized material specimens which are subjected to fluctuating stresses by a periodic loading. By a load (or load oscillation) we mean a continuous unimodal function on the unit interval, the value of which at any time gives the stress imposed by the deflection of the specimen. Let  $l_1, l_2, \dots$  be the sequence of loads which are to be applied at each oscillation so that at the  $i^{\text{th}}$  oscillation load  $l_i$  is imposed. We suppose that the loading is *cyclic* in the sense that for some  $m > 1$  and all  $i=1, \dots, m$

$$l_{j+m+1} = l_{k+m+1} \quad \text{for all } j \neq k \quad (2.1.1)$$

and the loading is continuous so that for all  $i=1, 2, \dots$

$$l_{i+1}(0) = l_i(1). \quad (2.1.2)$$

Hence the  $(j+1)^{\text{st}}$  cycle is the loading  $(l_{jm+1}, \dots, l_{j+m+1})$ .

We assume that fatigue failure is due to the initiation, growth and ultimate extension of a dominant crack. At each oscillation this crack is extended by some amount which is a random function due to the variation in the material, the magnitude of the imposed stress and a certain number of the prior loads and perhaps the actual crack extensions caused by the prior loads in that cycle.

Thus we now make our first assumption

- 1\* The incremental crack extension  $X_i$  following the application of the  $i^{\text{th}}$  oscillation is a random variable with a distribution which depends upon all the loads and actual crack extensions which have preceded it in that cycle.

The crack extension during the  $(j+1)^{\text{st}}$  cycle is

$$Y_{j+1} = X_{j+1} + \dots + X_{j+m} \quad \text{for } j=0,1,\dots,$$

where  $X_{j+1}$  is the (possibly microscopic) crack extension following the load  $z_i$  applied in the  $i^{\text{th}}$  oscillation of the  $(j+1)^{\text{st}}$  cycle.

It follows from Assumption 1\*, regardless of how much dependence exists between the successive random extensions per oscillation in each cycle, that the random total crack extensions per cycle are independent. Our notation will be

$$W_n = \sum_{j=1}^n Y_j,$$

which has distribution function

$$H_n(w) = P[W_n \leq w], \quad \text{for } n=1,2,\dots$$

It follows from elementary probability considerations, see p. 169, [6], that the distribution of  $C$ , the number of such cycles until failure, in the case failure is defined as the crack length exceeding some fixed critical length  $w$  for the first time, is

$$P[C \leq n] = 1 - H_n(w). \quad (2.2)$$

If we suppose that we are dealing with a long cycle of oscillations which is as complex as the ground-air-ground cycle in aeronautical failure studies, it might be reasonable to assume that each cycle itself consists of a large number of distinct phases of loading. Even though the total crack extension per cycle is the sum of random variables which are not necessarily independent or identically distributed, the number of summands

might be sufficiently large to make it reasonable to assume that the total crack extension per cycle is normally distributed.

Thus we could formally make a second assumption

2\* The total crack extension  $Y_j$  due to the  $j^{\text{th}}$  cycle is a normal random variable with mean  $\mu$  and variance  $\sigma^2$  for all  $j=1,2,\dots$

As a possible alternative one could postulate that the distribution of the crack extension might be different for the earlier cycles than it would be for the later ones. One of the first suppositions might be to make the distribution of the extension per cycle depend upon the size of the crack at the start of the cycle. One such assumption and the resulting distribution of the total crack length at the end of  $n$  cycles is presented in Section 4. We do not make that assumption here but instead, for reasons of simplicity, proceed with the study of the implications of the one above.

Note that the crack extension per cycle must be a nonnegative random variable. Thus in order for the Assumption 2\* to apply we must regard as negligible the probability with which this normal variable becomes negative e.g. by assuming

$$\mu > 3\sigma > 0. \quad (2.2.1)$$

It now follows from Assumptions 1\* and 2\* that the  $Y_j$  are independent and identically distributed normal random variables. Thus we have as the distribution of  $C$  from (2.2)

$$P[C \leq n] = 1 - \mathcal{N}\left(\frac{\omega - n\mu}{\sqrt{n}\sigma}\right) = \mathcal{N}\left(\frac{\sqrt{n}\mu}{\sigma} - \frac{\omega}{\sqrt{n}\sigma}\right) \quad (2.3)$$

where  $\mathcal{N}$  is the distribution function of the standard normal variate with zero mean and unit variance defined for  $-\infty < y < \infty$  by

$$\mathcal{N}(y) = \int_{-\infty}^y \frac{1}{\sqrt{2\pi}} e^{-t^2/2} dt. \quad (2.4)$$

We write

$$\alpha = \frac{\sigma}{\sqrt{\mu\omega}}, \quad \beta = \frac{\omega}{\mu} \quad (2.5)$$

and replace  $n$  by the nonnegative real variable  $t > 0$ . If we now denote the continuous extension of the discrete random variable  $C$  by  $T$ , a continuous nonnegative random variable, then it follows by (2.3) and (2.5) that  $T$  has the life distribution

$$F(t; \alpha, \beta) = \mathcal{N}\left[\frac{1}{\alpha} \xi(t/\beta)\right] \quad \text{for } t > 0 \quad (2.6)$$

where

$$\alpha > 0, \quad \beta > 0$$



and

$$f(t) = t^{\alpha} - t^{-\beta}, \quad (2.7)$$

This two-parameter family of distributions is a plausible model for the distribution of fatigue life. The set of all life distributions of the form (2.6) for  $\alpha, \beta > 0$  will be denoted by  $\mathcal{F}$  and most of this paper will be devoted to the study of its properties. We shall also refer to the law which has the distribution (2.6) with the somewhat shorter notation  $F(\alpha, \beta)$ .

Of course, there are motivational derivations for other distributions as well. The Weibull distribution, which is well known for its applications as a life length distribution and for fatigue life in particular, is obtained as a special case of the extreme value distributions, see p. 302 [7]. The Gamma family has also been obtained as a distribution of life by utilizing a model of a bundle of strands which are supporting a tensile load, see [4].

It is instructive to make a comparison between the derivation of the family of distributions  $\mathcal{F}$  and an appropriate adaptation of the classical heuristic argument, found for example p. 219, Cramér [5], as it might be used to obtain the lognormal distribution of the time until failure in fatigue. We present this derivation in Section 5.

### 3. SOME PROPERTIES OF $\mathcal{G}$

Let  $T$  have the distribution defined in (2.6).

Note that  $T$  is a two-parameter random variable with  $\beta$  as a location parameter since it is the median of the distribution. (We show later that  $\beta$  is neither the mode nor the mean.) Notice also that  $\alpha$  is a shape parameter and  $\beta$  a scale parameter. As we have seen  $\frac{1}{\alpha} \xi(T/\beta)$  is a standard normal random variable with mean zero and unit variance. If we let  $X$  be  $\mathcal{N}(0, \frac{\alpha^2}{4})$  we see that, in distribution,

$$2X = \xi(T/\beta). \quad (3.1)$$

If we define the function  $\psi$  by

$$\psi(x) = \xi^{-1}(2x) \text{ for all real } x, \quad (3.1.5)$$

then

$$T = \beta\psi(X) \quad (3.2)$$

From elementary algebra we find that

$$\psi(x) = [\rho(x)]^2 \quad (3.3)$$

where

$$\rho(x) = x + \sqrt{x^2 + 1}. \quad (3.4)$$

So by (3.2)

$$T = \beta[1 + 2X^2 + 2X\sqrt{1+X^2}] \quad (3.5)$$

where  $X$  is  $\mathcal{N}(0, \frac{\alpha^2}{4})$ . Hence we have immediately

$$E(T) = E(1 + \frac{\alpha^2}{2}) \quad (3.5.1)$$

$$\begin{aligned} E(T^2) &= \beta^2(1 + 2\alpha^2 + \frac{3\alpha^4}{2}) \\ \text{var}(T) &= (\alpha\beta)^2(1 + \frac{5\alpha^2}{4}) \end{aligned} \quad (3.5.2)$$

and we note that for fixed  $\alpha$  the variance of  $T$  increases as the scale parameter (median)  $\beta$  increases. This is not true for the lognormal distribution but empirical evidence shows that it is true for the fatigue lives themselves.

By noting that

$$\frac{1}{\rho(x)} = -x + \sqrt{1+x^2} = \rho(-x) \quad (3.5.3)$$

we see that, whenever  $-X$  has the same distribution as  $X$ , we have by (3.3), in distribution,

$$\frac{1}{\rho(X)} = \rho(X) \quad \text{and} \quad \frac{1}{\psi(X)} = \psi(X). \quad (3.6)$$

Thus there follows immediately from (3.2) the

Theorem 3.1. If  $T$  has a fatigue life distribution  $F(\alpha, \beta)$ , in  $\mathcal{F}$  then  $\frac{1}{T}$  has a distribution in  $\mathcal{F}$  given by  $F(\alpha, \frac{1}{\beta})$ . Moreover, for any real  $a > 0$ , the random variable  $aT$  has a distribution in  $\mathcal{F}$  given by  $F(\alpha, a\beta)$ .

It is known that every random variable with distribution defined by (2.2) for which the  $Y_j$  are nonnegative and have densities which are Pólya frequency functions of order 2, has an increasing failure rate, see [1]. We now show that the random variable  $T$  does not have this property and moreover the failure rate does not even increase on the average which is a weaker condition, see [2]. However, it just barely fails to have this last condition satisfied.

Remark 3.2.  $T$  does not have a failure rate which increases on the average.

Proof. Without loss of generality take  $\alpha = \beta = 1$  and define

$$Q(t) = -\ln(1 - \mathcal{N}[\xi(t)]).$$

Writing Mill's ratio

$$M(\xi) = \xi[1 - \mathcal{N}(\xi)]/\mathcal{N}'(\xi)$$

then applying L'Hospital's rule to  $Q(t)/t$  for  $t \rightarrow \infty$  and using (4.1.1) of the next section we obtain

$$\frac{Q(t)}{t} = \frac{\mathcal{N}'[\xi(t)]}{1 - \mathcal{N}[\xi(t)]} \xi'(t) = \frac{(1-t^{-2})}{2M[\xi(t)]}.$$

We know that as  $\xi \rightarrow \infty$  we have

$$M(\xi) \sim 1 - \frac{1}{\xi^2} + \frac{3}{\xi^4} + O(\xi^{-5}).$$

Hence  $\lim_{t \rightarrow \infty} \frac{Q(t)}{t} = \frac{1}{2}$  but  $Q(1) = \frac{1}{2} > \frac{1}{2}$ . This proves our contention

that  $\frac{Q(t)}{t}$  does not always increase. ||

Actual numerical calculation shows that  $\frac{Q(t)}{t}$  decreases slowly for  $t > 1.64$  as Figure 1 shows. Since normal random variables do have densities which are Pólya frequency functions of order 2 the reason the failure rate can decrease is that our summands fail to be nonnegative with probability one.

We do not regard this as being a serious shortcoming to our family of distributions when applied to fatigue, any more than the negativity of the normal distribution vitiates its usefulness when applied to such things as measurements, which theoretically cannot be negative.

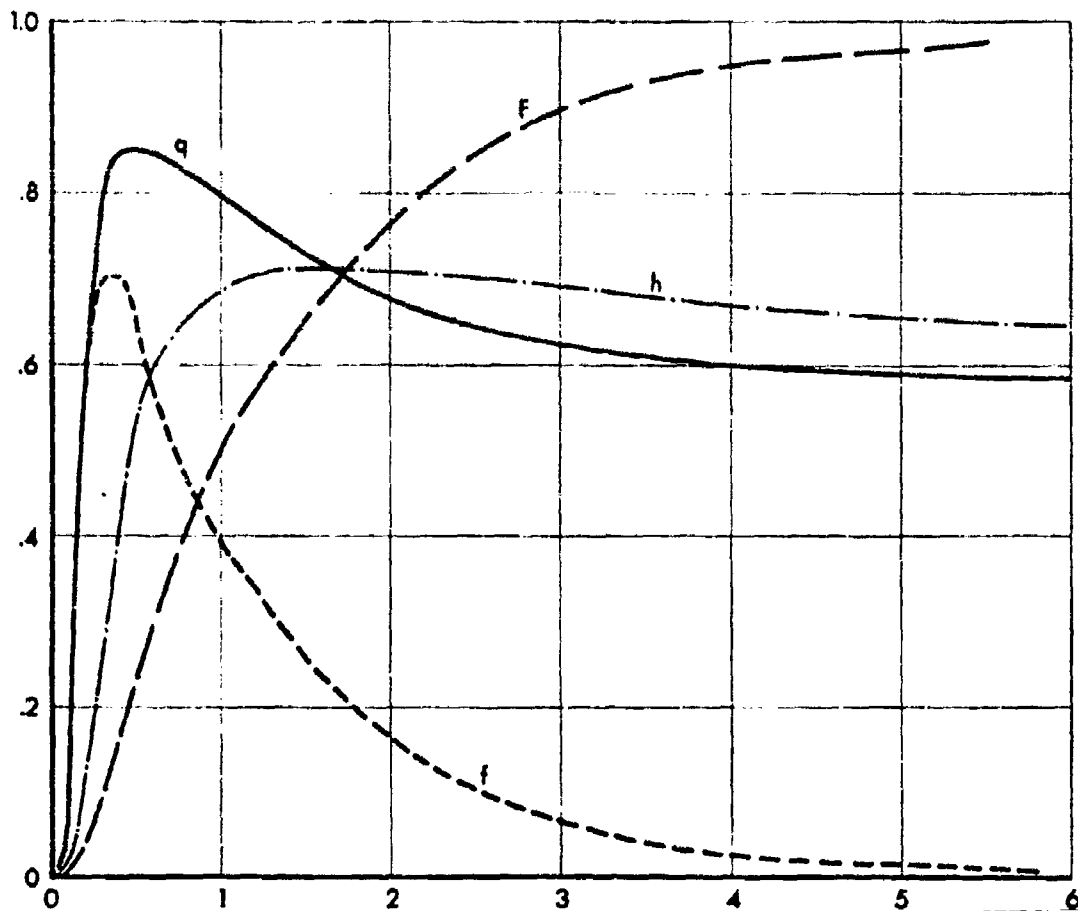


Figure 1

Graphs of the density  $f$ , the distribution  $F$ , the hazard rate  $q$ , and the average hazard rate  $h$  where for  $t > 0$

$$f(t) = \frac{t^{1/2} + t^{-1/2}}{2t\sqrt{2\pi}} \exp\left\{-\frac{t}{2} - \frac{1}{2t} + 1\right\}$$

$$F(t) = \int_0^t f(x)dx, \quad q(t) = \frac{f(t)}{1-F(t)}, \quad h(t) = \frac{1}{t} \int_0^t q(x)dx.$$

#### 4. A SECOND MODEL FOR THE DISTRIBUTION OF LIFE

In this section we replace the Assumption 2<sup>a</sup> by the more general one

- 2' The crack extension  $Y_{n+1}$  during the  $(n+1)^{\text{st}}$  cycle, given that the total crack length was  $s$  at the start of the cycle, is a normal random variable with mean  $\mu + \delta s$ , for some constant  $\delta \geq 0$ , and variance  $\sigma^2$  for each  $n=1,2,\dots$ .

By formula (2.2), to find the distribution of  $C$  it is necessary and sufficient that we find the distribution of  $W_n$ , called  $H_n$ .

By assumption

$$P[Y_{n+1} \leq y \mid W_n = s] = \mathcal{N}\left(\frac{y - \mu - \delta s}{\sigma}\right)$$

and for  $n=1$

$$H_1(y) = \mathcal{N}\left(\frac{y - \mu}{\sigma}\right) \quad -\infty < y < \infty. \quad (4.1)$$

Now by definition, setting  $y = x - s$

$$H_{n+1}(x) = \int_{-\infty}^{\infty} \mathcal{N}\left(\frac{y - \mu - \delta s}{\sigma}\right) dH_n(s). \quad (4.2)$$

We can now prove the

Theorem 4.1 The total crack length  $W_n$ , at the end of the  $n^{\text{th}}$  cycle, is normal with mean  $\mu_n$  and variance  $\sigma_n^2$  where

$$\mu_n = \frac{\mu}{\delta} [(1+\delta)^n - 1], \quad \sigma_n^2 = \sigma^2 \frac{(1+\delta)^{2n} - 1}{(1+\delta)^2 - 1}. \quad (4.3)$$

Proof by induction. The statement is true for  $n=1$ . Assume it true for  $n$ , then from (4.2)

$$H_{n+1}(x) = \int_{-\infty}^{\infty} \mathcal{N}\left(\frac{u-s}{v}\right) d_s \mathcal{N}\left(\frac{s-\mu_n}{\sigma_n}\right)$$

where

$$u = \frac{x-\mu}{1+\delta} \quad v = \frac{\sigma}{1+\delta}.$$

But this we recognize as the usual convolution of two normal random variables and hence the result is

$$H_{n+1}(x) = \mathcal{N}\left(\frac{u-\mu_n}{\sqrt{v^2+\sigma_n^2}}\right).$$

Hence by simplification we find that

$$\mu_{n+1} = \mu + (1+\delta)\mu_n, \quad \sigma_{n+1}^2 = \sigma^2 + (1+\delta)^2\sigma_n^2$$

and one checks that formulas given do satisfy the recursion relations (4.3). ||

Strictly speaking a proper distribution, analogous to that obtained in (2.6), which would be the continuous extension of the present case, cannot be generated since  $P[C < \infty] < 1$ .

For, one can see that

$$P[C \leq n] = \mathcal{N}\left(\frac{\mu_n - \omega}{\sigma_n}\right)$$

and as  $n \rightarrow \infty$  we have  $\mu_n/\sigma_n \rightarrow \frac{\mu}{\delta\sigma} \sqrt{(1+\delta)^2 - 1} < \infty$ .

Of course this would not be of practical significance, since we are restricting ourselves to situations where  $\frac{\mu}{\sigma} > 3$ , see (2.2.1), hence  $\lim_{n \rightarrow \infty} \frac{\mu_n}{\sigma_n} > 3\sqrt{2/\delta}$  for  $0 < \delta < 1$ , and  $\lim_{n \rightarrow \infty} \frac{\mu_n}{\sigma_n} > 3$  for any  $\delta$ .

## 5. A DERIVATION OF THE LOGNORMAL DISTRIBUTION

In our notation the  $i^{\text{th}}$  oscillation is caused by the loading function  $l_i$ , and the crack extension depends upon both the loading function and the size of the crack already attained at the time when the oscillation was begun.

Consider the  $n$  loads  $l_1, \dots, l_n$  which have resulted in a crack of size  $W_n$ , then

$$W_{n+1} - W_n = \frac{L_{n+1}}{\delta(W_n)},$$

assuming the incremental growth of the crack is proportional to some random functional  $L_{n+1}$  of the loading function  $l_{n+1}$  imposed at that time, and some function of the size of the crack  $\delta(W_n)$ . Now we have

$$L_1 + \dots + L_n = \sum_{j=1}^n (W_j - W_{j-1}) \delta(W_{j-1}).$$

Letting  $n$  become large we then have the left-hand side becoming asymptotically a normal random variable, due to the periodicity of loading, while the right-hand side becomes in the limit

$$\int_{Z_0}^Z \delta(t) dt.$$

If we take  $\delta(t) = 1/t$  then  $Z$  becomes asymptotically lognormal.

The important distinction between this derivation and that which we previously considered is that while a given number of identically distributed summands does approach normality as the number gets large, a random number with large mean of identically distributed summands does not necessarily do so. Moreover, we maintain that what is at question here is the number of impulses required to exceed a certain critical crack size and it is indeed a random number of cycles which will be necessary to accomplish this.



## 6. CONCLUSION

A derivation, based on plausible physical considerations, for a family of distributions is, by itself, not a conclusive argument that such a particular family should be used in life studies. No family, however reasonable its derivation, can be accepted for use in fatigue life studies until it is confronted with actual fatigue data obtained under various conditions and the distribution is shown to represent adequately the life length  $\Delta$  which are obtained.

In order to do this one must have the theory of estimation for this family completed. The derivation of parametric estimators and the ancillary computing formulas for this family will be presented in a later study. Also further studies of the application of this distribution to the calculation of "safe life" will be made. Thus the confrontations of this family with actual data will be carried out, to provide the justification for this presentation.

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